



Spin squeezing on the caesium clock transition: Interferometric quantum non-demolition measurement on cold atoms

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Experimentalists

Current

Ulrich Busck Hoff
Patrick Windpassinger
Daniel Oblak
Jürgen Appel
Niels Kjærgaard
Eugene Polzik

Previous







Mark Saffman
Carlos Garrido Alzar
Plamen Petrov
Jens Lykke Sørensen
Jens Mikkelsen
Anton Vershovski
Wolfgang Tittel





Overview of presentation

This presentation will contain:

- Motivation and applications 
- Atomic system, interferometer and quantum variables 
- QND interaction 
- Experimental setup and characterisation 
- Measurements of atomic noise 
- Things to come (soon) 

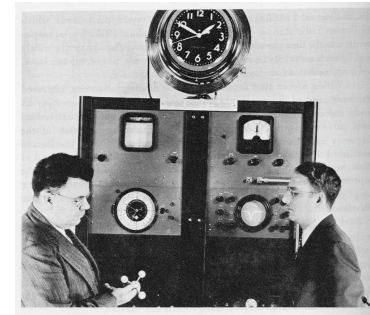


Motivation and applications

What's it all good for?

Increased sensitivity in spectroscopy^{1,2,3}

Atomic magnetometers
Microwave fountain clocks
Optical lattice clocks



Quantum memory⁴

Increased fidelity

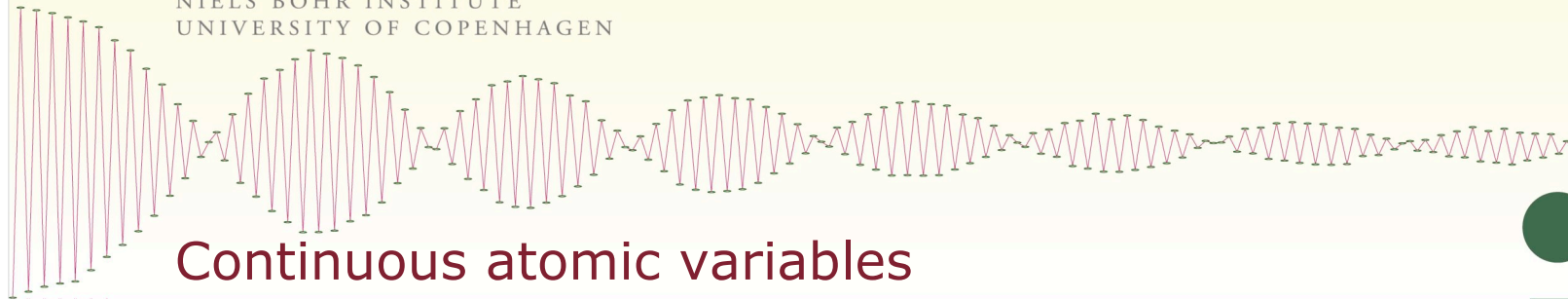


Non classical states for quantum information

Qubit system



1. Auzinsh et al. PRL **93** (17) 173002
2. Wineland, Bollinger, Itano, and Heinzen. PRA **50** (1) 67
3. Andre, Sørensen, and Lukin. PRL **92**, 230801
4. Hammerer, Mølmer, Polzik, and Cirac. PRA **70** 044304



Continuous atomic variables

Pseudo-spin on the Bloch sphere^{1,2}



Atomic pseudo-spin

$$\hat{j}_x^{(k)} = \frac{1}{2}(\hat{\sigma}_{34}^{(k)} + \hat{\sigma}_{43}^{(k)})$$

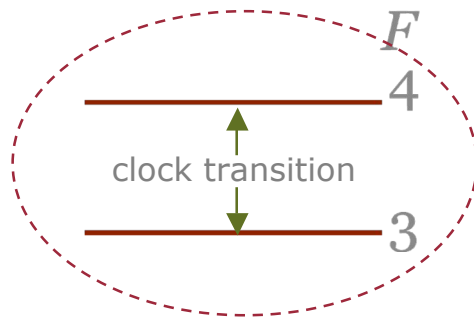
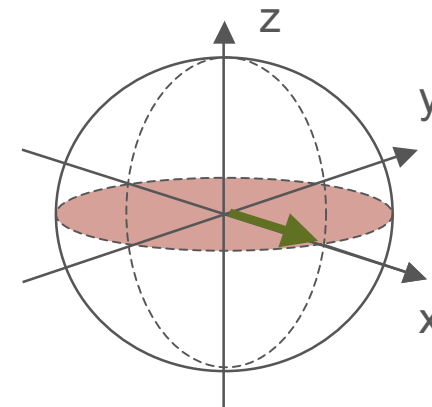
$$\hat{j}_y^{(k)} = \frac{-i}{2}(\hat{\sigma}_{34}^{(k)} - \hat{\sigma}_{43}^{(k)})$$

$$\hat{j}_z^{(k)} = \frac{1}{2}(\hat{\sigma}_{33}^{(k)} - \hat{\sigma}_{44}^{(k)})$$

Coherent spin state

$$|\Psi\rangle_k = \frac{1}{\sqrt{2}}|3\rangle_k + \frac{1}{\sqrt{2}}|4\rangle_k$$

Bloch sphere - qubit

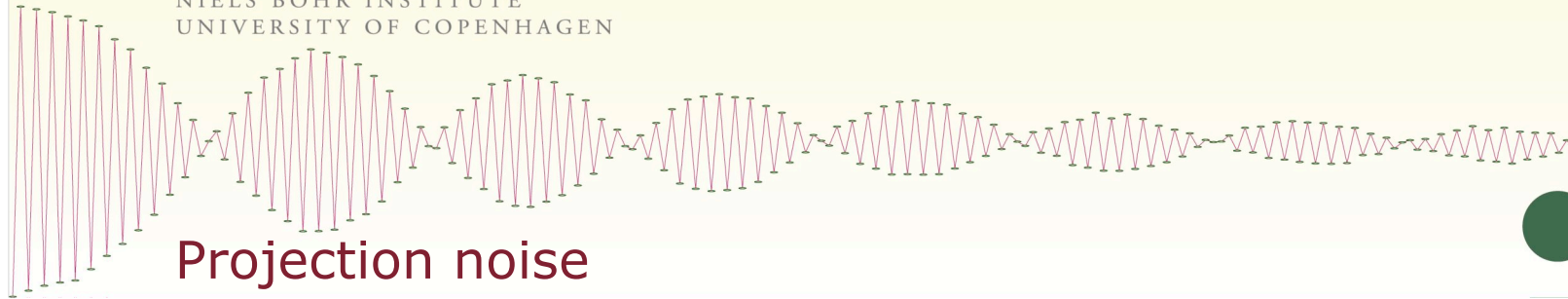


Collective operators

$$\hat{J}_i = \sum_{k \in V} \hat{j}_i^{(k)}$$

$$\hat{N} = \sum_{k \in V} \hat{\sigma}_{33}^{(k)} + \sum_{k \in V} \hat{\sigma}_{44}^{(k)}$$

1. Wineland, Bollinger, Itano, and Heinzen. PRA **50** (1) p.67
2. Saffman Oblak, Appel, & Polzik et. al. preprint available from speaker



Projection noise

Expectation values^{1,2}



Mean values

$$\langle \hat{J}_x \rangle = \frac{1}{2} \langle \hat{N} \rangle$$

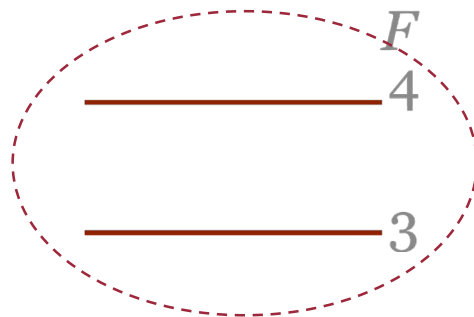
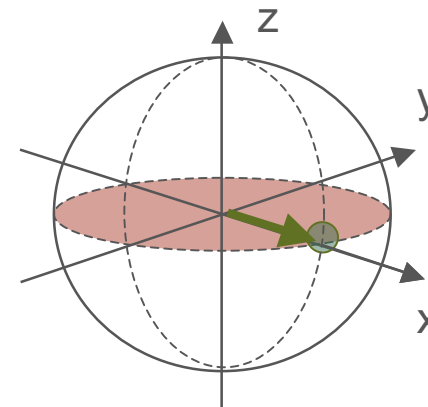
$$\langle \hat{J}_y \rangle = 0$$

$$\langle \hat{J}_z \rangle = 0$$

Coherent spin state

$$|\Psi\rangle_k = \frac{1}{\sqrt{2}} |3\rangle_k + \frac{1}{\sqrt{2}} |4\rangle_k$$

Bloch sphere - qubit



Uncertainties

$$\langle (\delta \hat{J}_x)^2 \rangle = 0$$

$$\langle (\delta \hat{J}_y)^2 \rangle = \frac{1}{2} \langle \hat{J}_x \rangle = \frac{1}{4} \langle \hat{N} \rangle$$

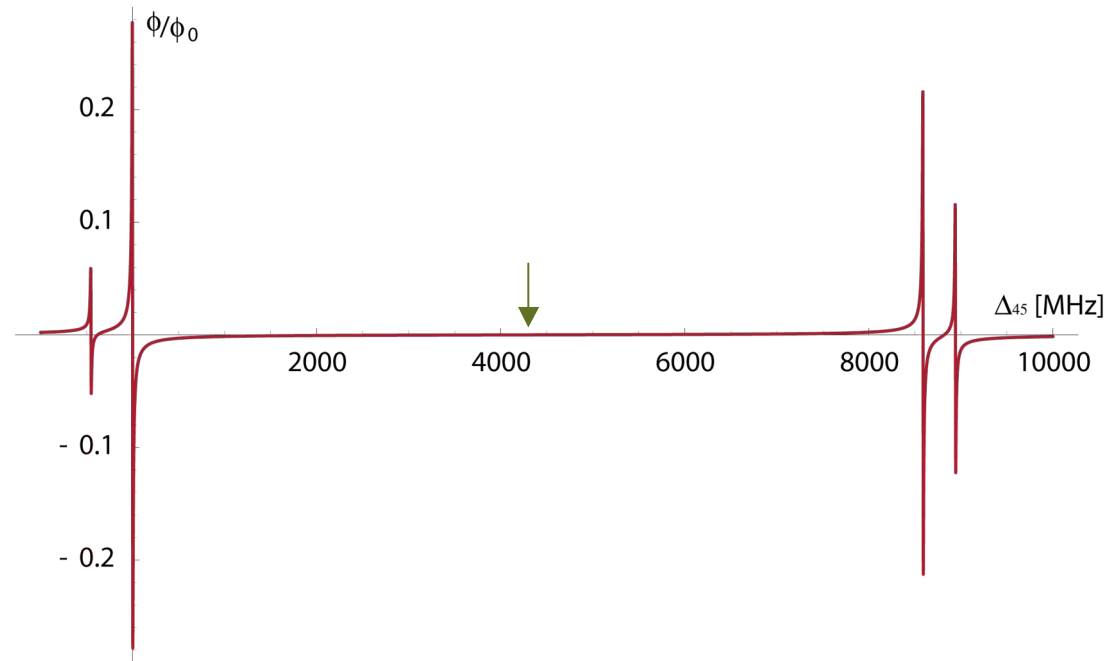
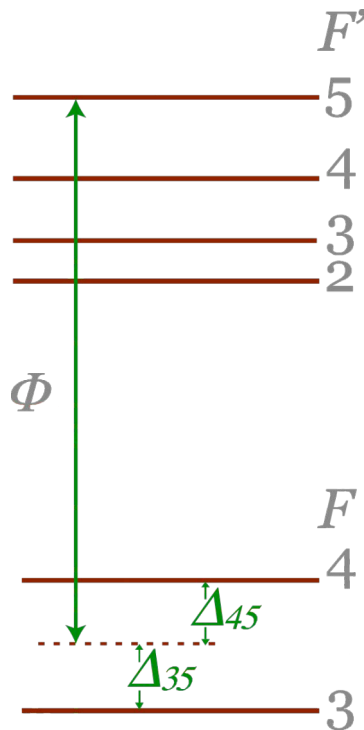
$$\langle (\delta \hat{J}_z)^2 \rangle = \frac{1}{2} \langle \hat{J}_x \rangle = \frac{1}{4} \langle \hat{N} \rangle$$

1. Wineland, Bollinger, Itano, and Heinzen. PRA **50** (1) p.67
2. Saffman Oblak, Appel, & Polzik et. al. preprint available from speaker



Dispersive interaction with light

Phase-shift of probe beams



$$\phi = \phi_0 \left[\sum_{e=2,4} c_{3e} \frac{\Delta_{3e} \frac{\gamma}{2}}{\Delta_{3e}^2 + (\frac{\gamma}{2})^2} \hat{N}_3 + \sum_{e=3,5} c_{4e} \frac{\Delta_{4e} \frac{\gamma}{2}}{\Delta_{4e}^2 + (\frac{\gamma}{2})^2} \hat{N}_4 \right]$$

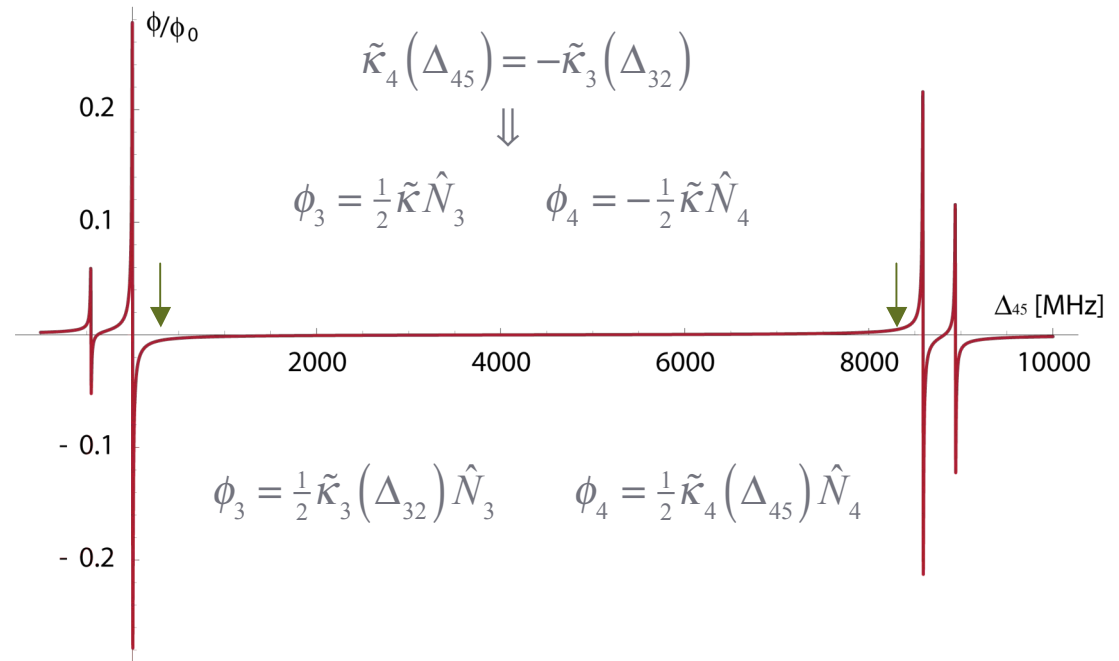
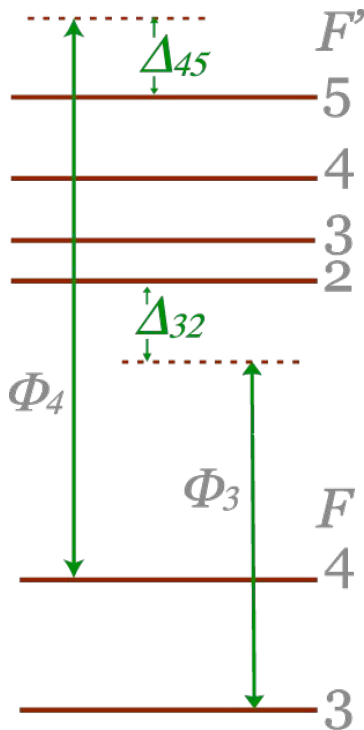
$$= \frac{1}{2} \left[\tilde{\kappa}_3 (\Delta_{32}) \hat{N}_3 + \tilde{\kappa}_4 (\Delta_{45}) \hat{N}_4 \right]$$

c_{3e} : transition strength, Δ_{3e} : detuning, γ : natural linewidth



Dispersive interaction with light

Phase-shift of probe beams



$$\phi = \phi_0 \left[\sum_{e=2,4} c_{3e} \frac{\Delta_{3e} \frac{\gamma}{2}}{\Delta_{3e}^2 + (\frac{\gamma}{2})^2} \hat{N}_3 + \sum_{e=3,5} c_{4e} \frac{\Delta_{4e} \frac{\gamma}{2}}{\Delta_{4e}^2 + (\frac{\gamma}{2})^2} \hat{N}_4 \right]$$

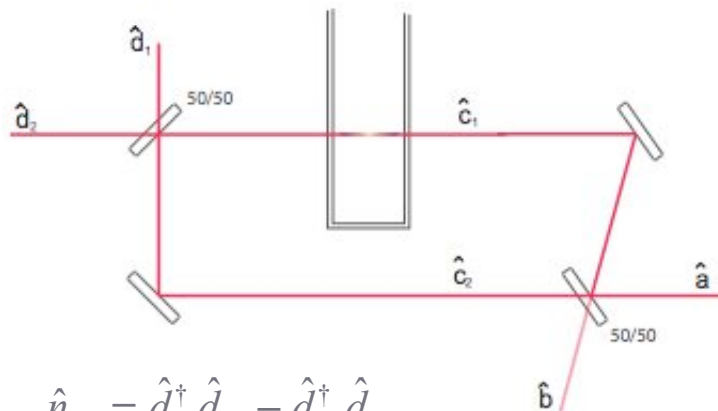
$$= \frac{1}{2} \left[\tilde{\kappa}_3(\Delta_{32}) \hat{N}_3 + \tilde{\kappa}_4(\Delta_{45}) \hat{N}_4 \right]$$

1. Petrov, Oblak et al. PRA 75, 033803 (2007)



Interferometric measurement

An intuitive picture



$$\tilde{\kappa}_4(\Delta_{45}) = -\tilde{\kappa}_3(\Delta_{32})$$

⇓

$$\phi_3 = \frac{1}{2}\tilde{\kappa}\hat{N}_3 \quad \phi_4 = -\frac{1}{2}\tilde{\kappa}\hat{N}_4$$

$$\hat{n}_{k-} = \hat{d}_{k1}^\dagger \hat{d}_{k1} - \hat{d}_{k2}^\dagger \hat{d}_{k2}$$

$$i_- \propto \langle \hat{n}_- \rangle \equiv \langle \hat{n}_{3-} + \hat{n}_{4-} \rangle$$

$$= \langle \hat{n}_3 \rangle \cos(\phi_3 + \phi') + \langle \hat{n}_4 \rangle \cos(\phi_4 + \phi')$$

$$\approx 2\tilde{\kappa}_3 \langle \hat{n}_3 \rangle \hat{N}_{3z} - 2\tilde{\kappa}_4 \langle \hat{n}_4 \rangle \hat{N}_{4z}$$

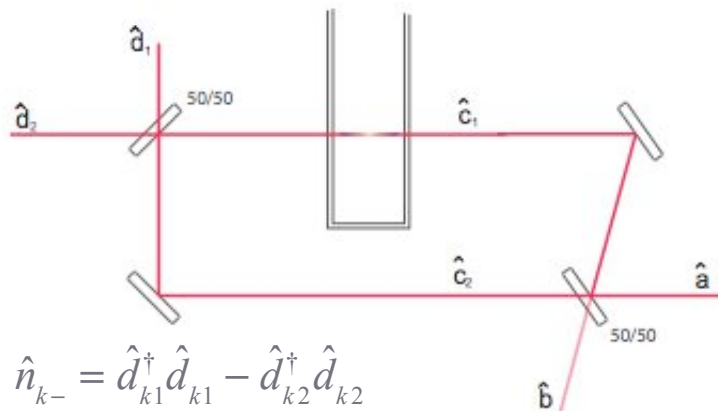
$$= 2\tilde{\kappa} \langle \hat{n} \rangle [\hat{N}_{3z} - \hat{N}_{4z}] = 2\tilde{\kappa}\tau \langle \hat{n} \rangle \hat{J}_z$$

interferometer optical
path-length offset: $\pi/2$

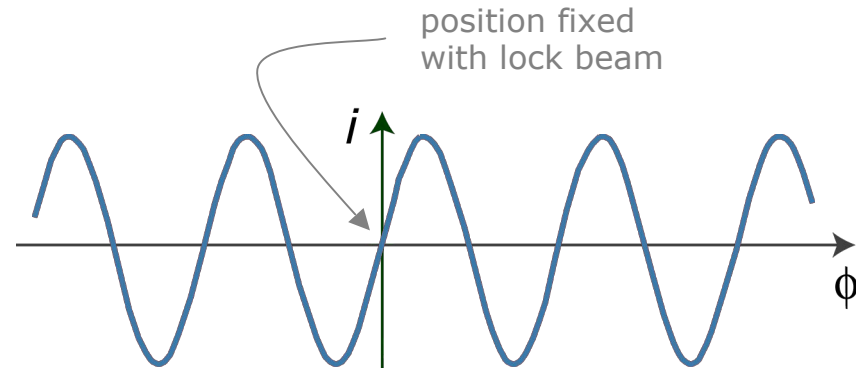


Interferometric measurement

An intuitive picture

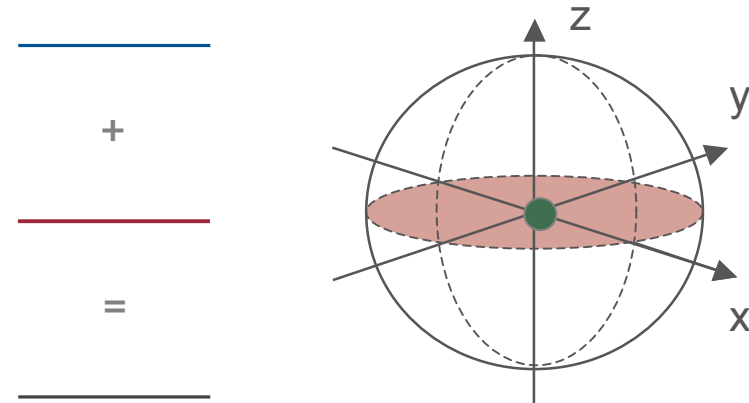


$$\begin{aligned} i_- &\propto \langle \hat{n}_- \rangle \equiv \langle \hat{n}_{3-} + \hat{n}_{4-} \rangle \\ &= \langle \hat{n}_3 \rangle \cos(\phi_3 + \phi') + \langle \hat{n}_4 \rangle \cos(\phi_4 + \phi') \\ &\approx 2\tilde{\kappa}_3 \langle \hat{n}_3 \rangle \hat{N}_{3z} + 2\tilde{\kappa}_4 \langle \hat{n}_4 \rangle \hat{N}_{4z} \\ &= 2\tilde{\kappa} \langle \hat{n} \rangle [\hat{N}_{3z} - \hat{N}_{4z}] = 2\tilde{\kappa}\tau \langle \hat{n} \rangle \hat{J}_z \end{aligned}$$



Output signal

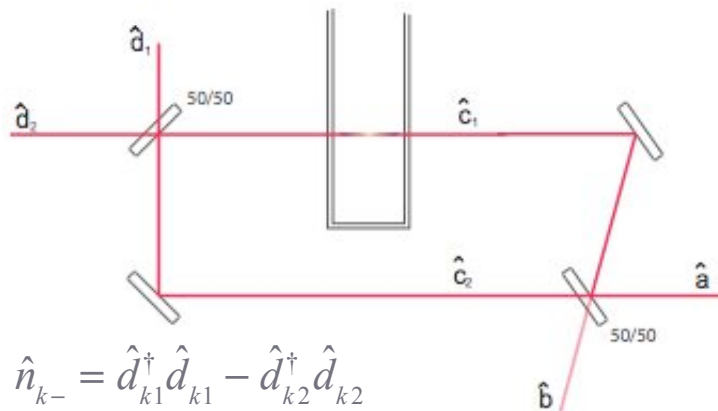
Atomic state





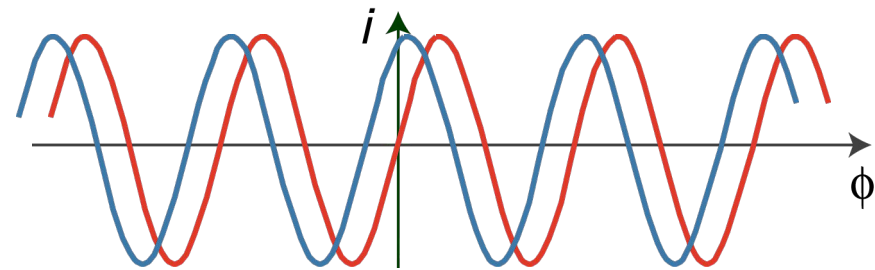
Interferometric measurement

An intuitive picture



$$\hat{n}_{k-} = \hat{d}_{k1}^\dagger \hat{d}_{k1} - \hat{d}_{k2}^\dagger \hat{d}_{k2}$$

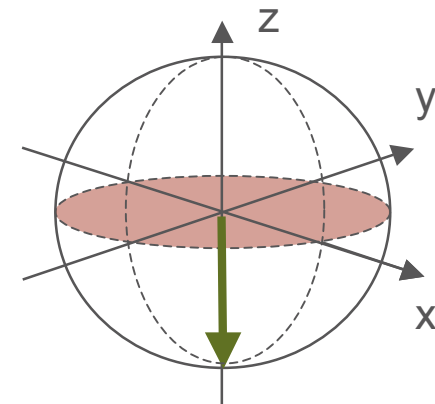
$$\begin{aligned} i_- &\propto \langle \hat{n}_- \rangle \equiv \langle \hat{n}_{3-} + \hat{n}_{4-} \rangle \\ &= \langle \hat{n}_3 \rangle \cos(\phi_3 + \phi') + \langle \hat{n}_4 \rangle \cos(\phi_4 + \phi') \\ &\approx 2\tilde{\kappa}_3 \langle \hat{n}_3 \rangle \hat{N}_{3z} + 2\tilde{\kappa}_4 \langle \hat{n}_4 \rangle \hat{N}_{4z} \\ &= 2\tilde{\kappa} \langle \hat{n} \rangle [\hat{N}_{3z} - \hat{N}_{4z}] = 2\tilde{\kappa}\tau \langle \hat{n} \rangle \hat{J}_z \end{aligned}$$



+



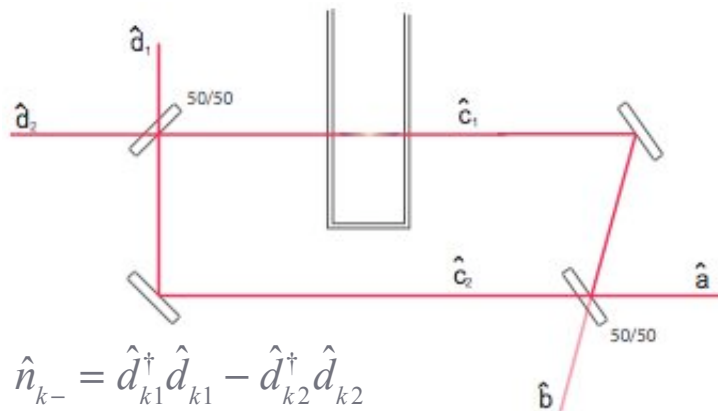
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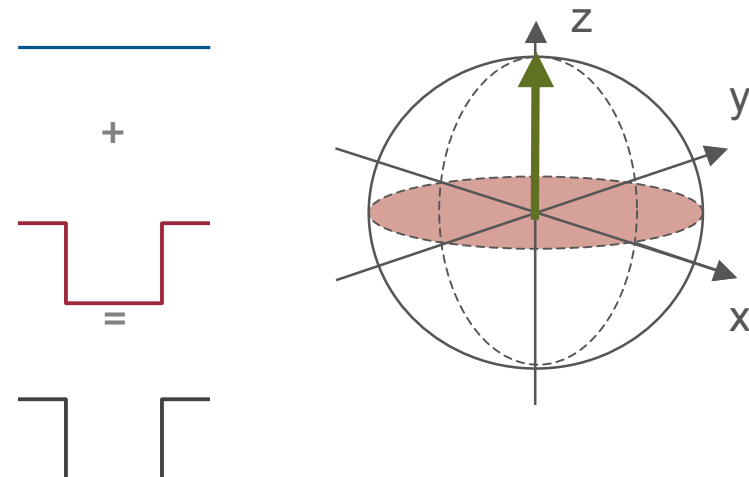
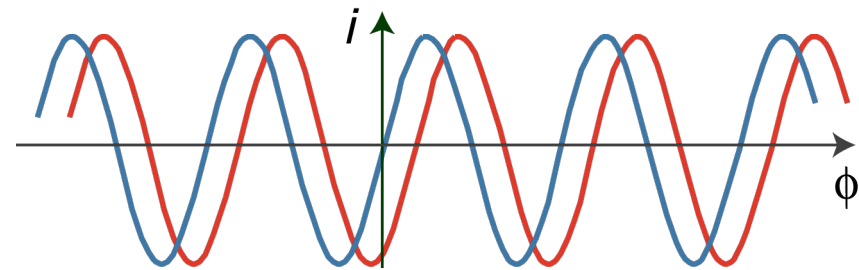
Interferometric measurement

An intuitive picture



$$\hat{n}_{k-} = \hat{d}_{k1}^\dagger \hat{d}_{k1} - \hat{d}_{k2}^\dagger \hat{d}_{k2}$$

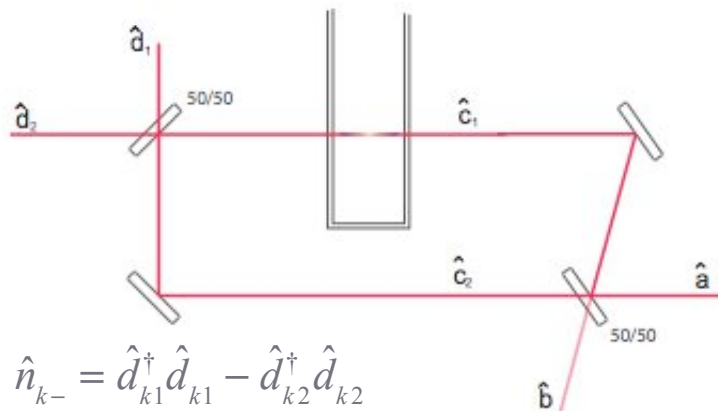
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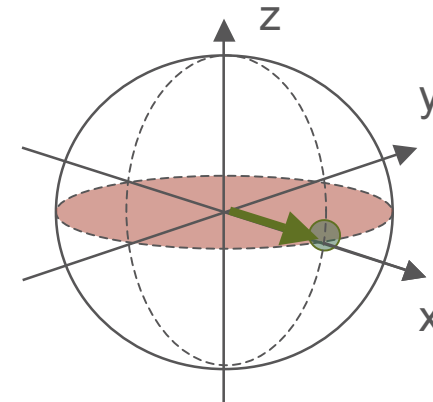
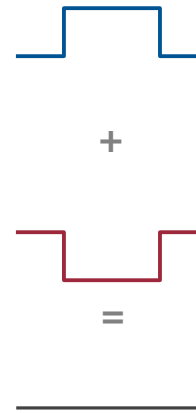
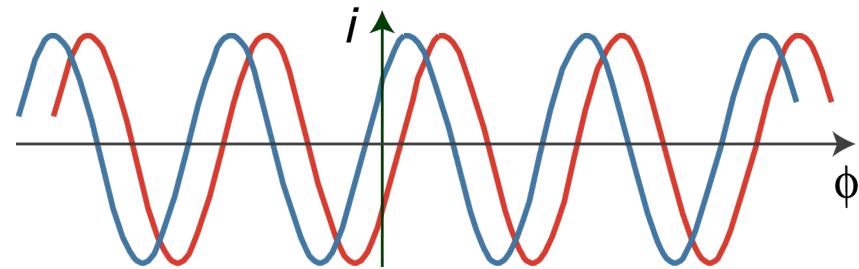
Interferometric measurement

An intuitive picture



$$\hat{n}_{k-} = \hat{d}_{k1}^\dagger \hat{d}_{k1} - \hat{d}_{k2}^\dagger \hat{d}_{k2}$$

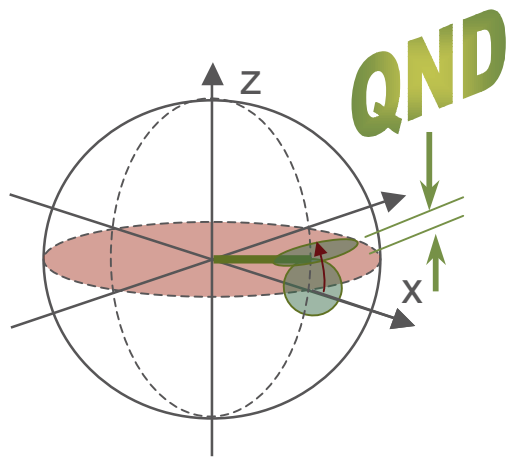
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Quantum non-demolition measurement

Squeezing by interaction and subsequent measurement of light pulse



Output variances

$$(\delta \hat{J}_y^{out})^2 = (\delta \hat{J}_y^{in})^2 + \tilde{\kappa} \tau \langle \hat{N} \rangle (\delta \hat{n}_\perp^{in})^2 = \frac{1}{4} \langle \hat{N} \rangle (1 + 4\tilde{\kappa} \tau \langle \hat{n} \rangle)$$

$$(\delta \hat{J}_y^{out})^2 (\delta \hat{J}_z^{out})^2 = \frac{1}{16} \langle \hat{N} \rangle^2 \Rightarrow (\delta \hat{J}_z^{out})^2 = \frac{\frac{1}{4} \langle \hat{N} \rangle}{(1 + 4\tilde{\kappa} \tau \langle \hat{n} \rangle)}$$

QND interaction Hamiltonian

$$\hat{H}_{int} = \hbar 4\tilde{\kappa} \frac{1}{2} (\hat{n} + \hat{n}_\perp) \hat{J}_z \Rightarrow$$

$$\hat{n}_{k\perp} = -i(\hat{d}_{k1}^\dagger \hat{d}_{k2} - \hat{d}_{k2}^\dagger \hat{d}_{k1}), \quad \hat{n}_\perp = \hat{n}_{3\perp} + \hat{n}_{4\perp}$$

$$\hat{n}_k = \hat{d}_{k1}^\dagger \hat{d}_{k1} + \hat{d}_{k2}^\dagger \hat{d}_{k2}, \quad \hat{n} = \hat{n}_3 + \hat{n}_4$$

CEWQO 2008 Belgrade, Serbia

Operator input-output relations

$$n_\perp^{out} = n_\perp^{in}, \quad n_-^{out} = n_-^{in} + 2\tilde{\kappa} \tau \langle \hat{n} \rangle \hat{J}_z$$

$$\hat{J}_z^{out} = \hat{J}_z^{in}, \quad \hat{J}_y^{out} = \hat{J}_y^{in} + 2\tilde{\kappa} \tau \langle \hat{N} \rangle \frac{1}{2} (\langle \hat{n} \rangle + n_\perp^{in})$$

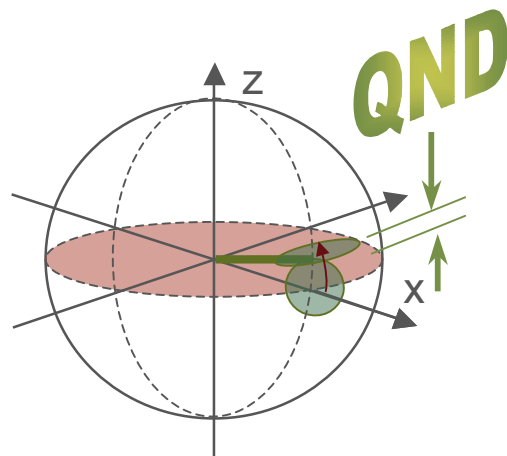
light-shift

$$\langle \hat{n}_\perp \rangle = 0, \quad (\delta \hat{n}_\perp)^2 = \langle \hat{n} \rangle, \quad \text{back-action}$$



Quantum non-demolition measurement

Squeezing by interaction and subsequent measurement of light pulse¹



Output variances²

$$(\delta \hat{J}_y^{out})^2 = (\delta \hat{J}_y^{in})^2 + \tilde{\kappa} \tau \langle \hat{N} \rangle (\delta \hat{n}_\perp^{in})^2 = \frac{1}{4} \langle \hat{N} \rangle (1 + 4\tilde{\kappa} \tau \langle \hat{n} \rangle)$$

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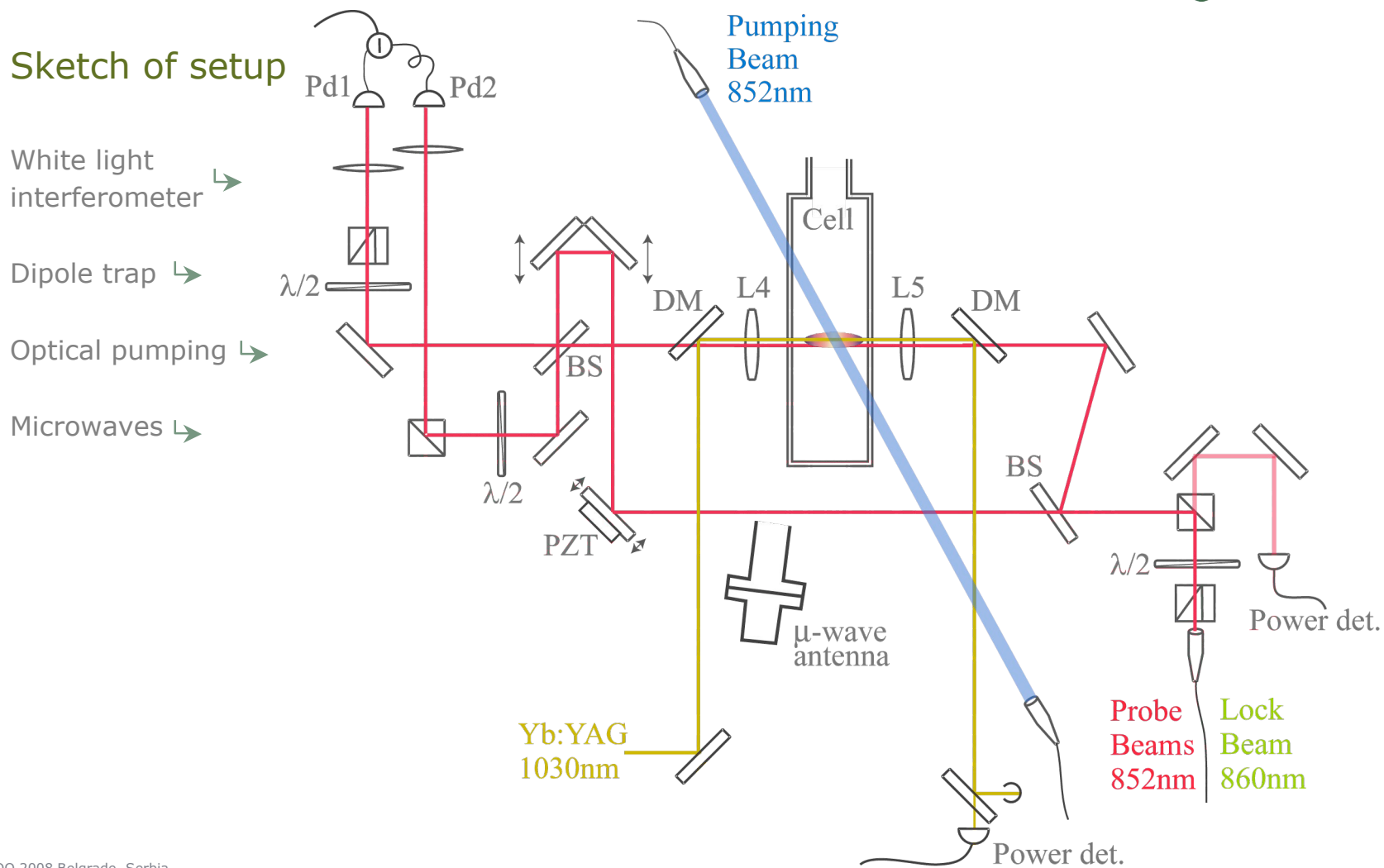
Squeezing criterion^{3,4}

$$(\delta \hat{J}_z^{out})^2 < \frac{1}{2} \langle \hat{J}_x^{out} \rangle = \frac{1}{4} \eta \langle \hat{N} \rangle$$

1. Kuzmich, Bigelow, and Mandel. *Eur. Phys. Lett.* **42** 481
2. Oblak et. al. *PRA* **71** 043807
3. Wineland, Bollinger, Itano, and Heinzen. *PRA* **50** (1) 67
4. Kitagawa and Ueda. *PRA* **47** 5138



Setup





Setup

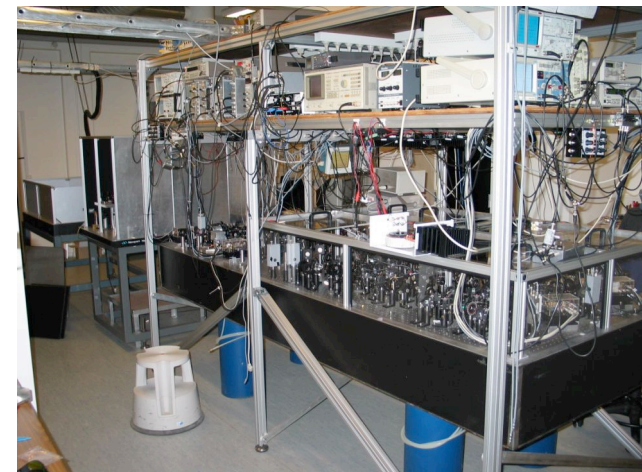
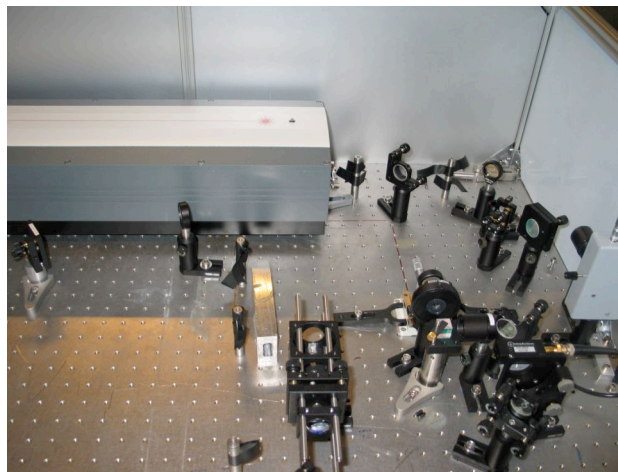
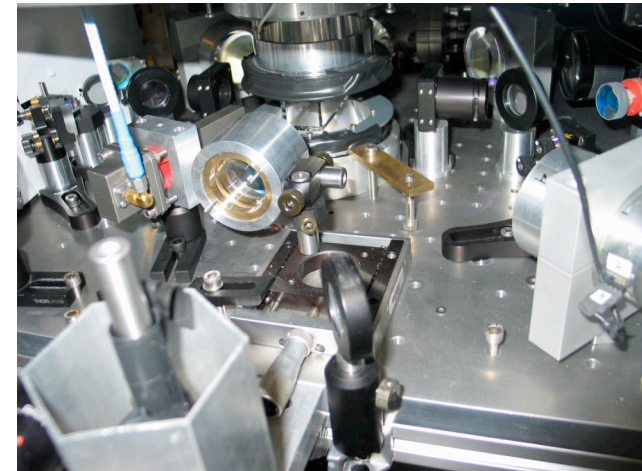
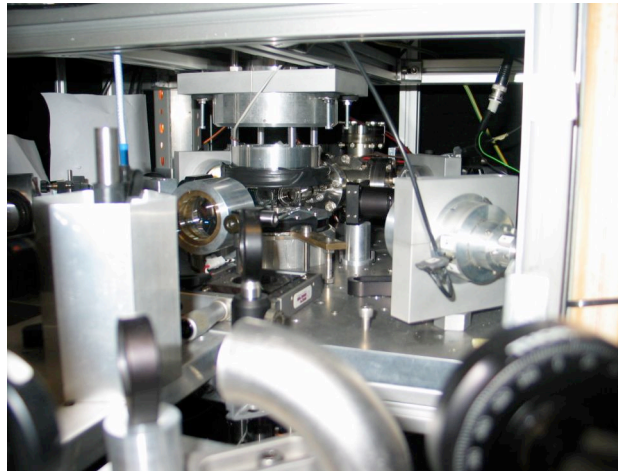
Photos of setup

White light
interferometer

Dipole trap

Optical pumping

Microwaves





Measurement procedure

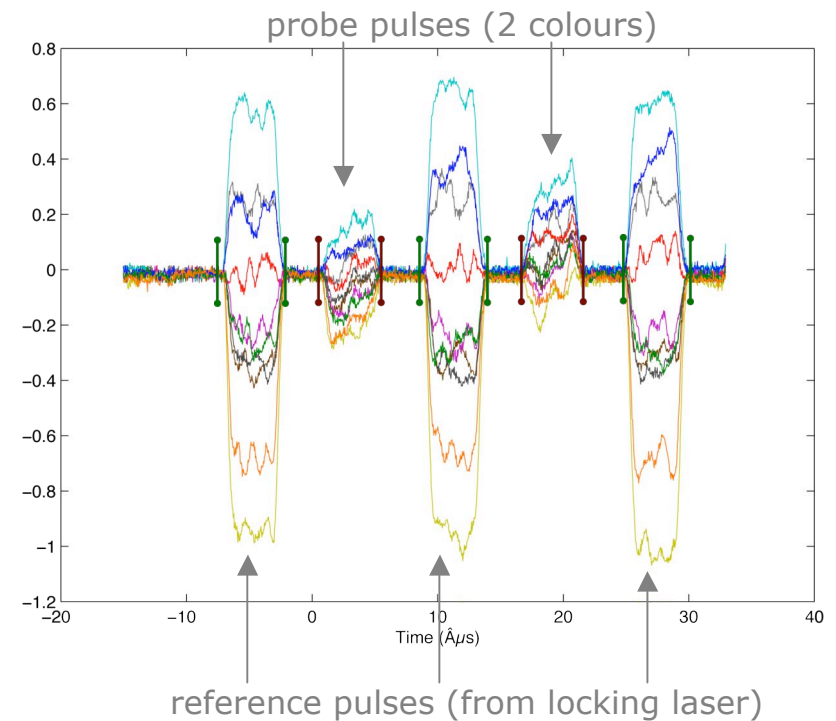
Pulsed measurement¹

4 μ s pulses generated in AOM

Incident on balanced low-noise detector

Detector output recorded on digital scope

Integration over pulses (window) and
further processing on PC





Measurement procedure

Pulsed measurement - corrections

Baseline subtraction & probe power normalisation

$$p_{j,p}^{(i)} = \rho^{(i)} \left(p_{j,p}^{(i)} - p_{j,bg}^{(i)} \right)$$

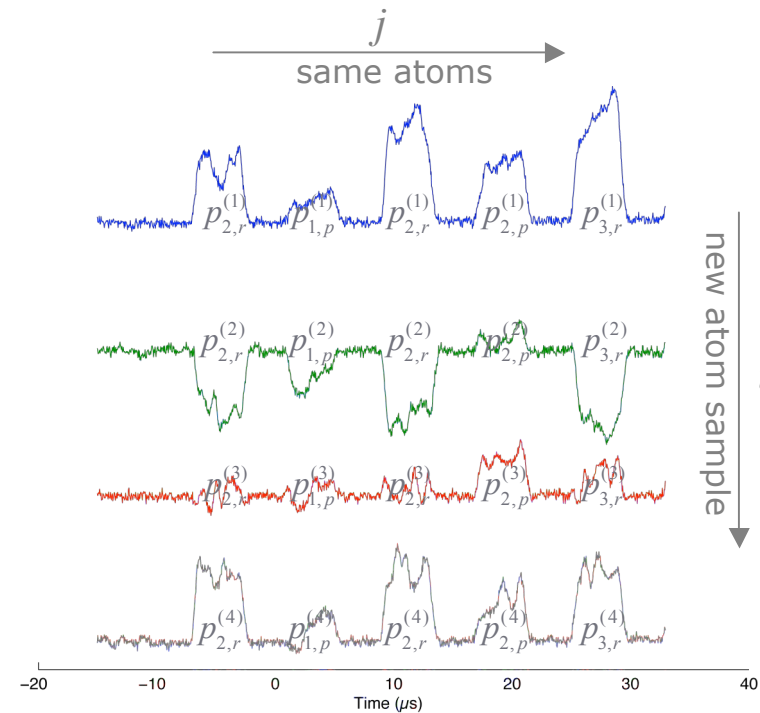
$$\rho^{(i)} = \sum_j^M P_{j,p}^{(i)} / \left(\frac{1}{KM} \sum_{i,j}^{K,M} P_{j,p}^{(i)} \right) \leftarrow \text{factor for power normalisation}$$

Reference pulse correction

$$p_{p,j}^{(i)} = p_{p,j}^{(i)} - \chi_{ref} \frac{1}{2} \left(p_{r,j}^{(i)} + p_{r,j+1}^{(i)} \right),$$

$$\chi_{ref} = \min \left\{ \frac{1}{K} \sum_{i=1}^K \left[p_{p,j}^{(i)} - \chi \frac{1}{2} \left(p_{r,j}^{(i)} + p_{r,j+1}^{(i)} \right) \right]^2 \right\}$$

\uparrow correlation factor between probe and reference pulses
 - classical fluctuations $\Rightarrow \chi_{ref} \approx 1$
 - shot noise $\Rightarrow \chi_{ref} \approx 0$





Measurement procedure

Pulsed measurement - noise calculation

Projection noise measurement

$$\left(\delta p_1\right)_{proj}^2 = \frac{1}{k-1} \sum_{i=1}^K \left[p_{p,1}^{(i)} \right]^2 \propto \frac{1}{4} \langle \hat{N} \rangle$$

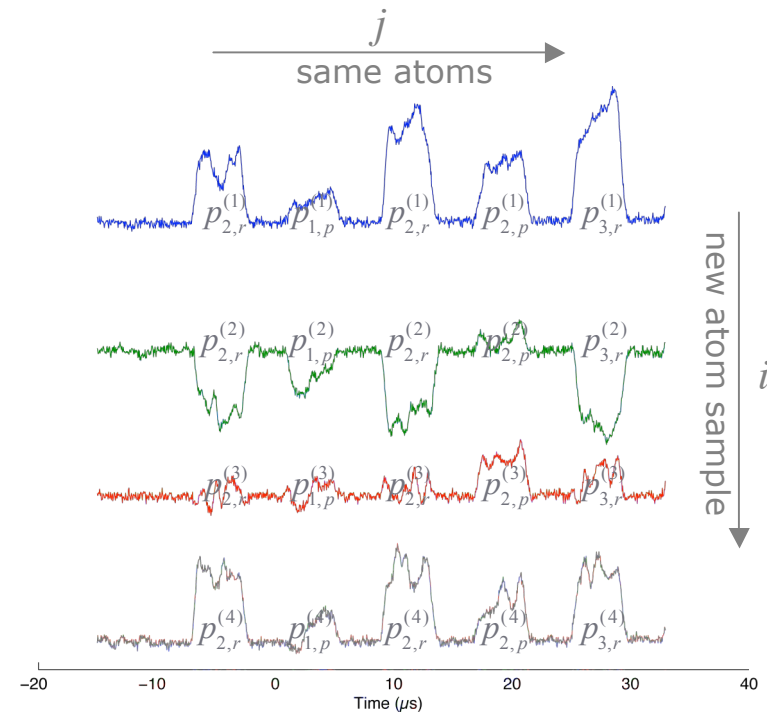
Squeezing measurement - in principle

$$\left(\delta p_{12}\right)_{sq}^2 = \frac{1}{2k-1} \sum_{i=1}^K \left[p_{p,2}^{(i)} - \chi_{sq} p_{p,1}^{(i)} \right]^2$$

$$\chi_{sq} = \min_{\chi} \left\{ \frac{1}{2k-1} \sum_{i=1}^K \left[p_{p,2}^{(i)} - \chi p_{p,1}^{(i)} \right]^2 \right\}$$

Experimental squeezing criterion

$$\left(\delta p_{12}\right)_{sq}^2 < \frac{1}{4} \eta \langle \hat{N} \rangle = \eta \left(\delta p_1\right)_{proj}^2$$



Pulse durations: 4 μs

Pulse separation: 16 μs

Probe power: 10 μW ~ 1.6 · 10⁸ photons

Reference power: 50 μW ~ 8 · 10⁸ photons



Near spin-squeezing

Experimental parameters

Projection noise measurement

$$(\delta p_1)_{proj}^2 = \frac{1}{k-1} \sum_{i=1}^K [p_{p,1}^{(i)}]^2 \propto \frac{1}{4} \langle \hat{N} \rangle$$

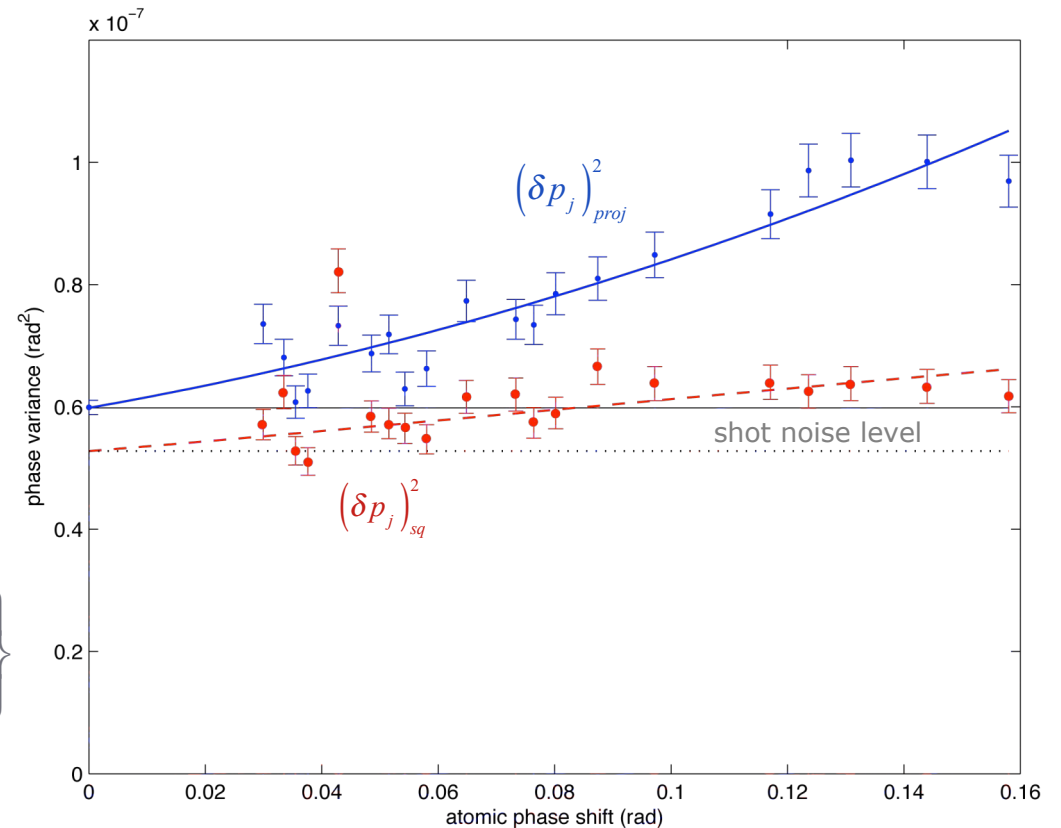
Squeezing measurement - in principle

$$(\delta p_j)_{sq}^2 = \frac{1}{2k-1} \sum_{i=1}^K [p_{p,j}^{(i)} - \chi_{sq} p_{p,j+1}^{(i)}]^2$$

$$\chi_{sq} = \min_{\chi} \left\{ \frac{1}{2k-1} \sum_{i=1}^K [p_{p,j}^{(i)} - \chi p_{p,j+1}^{(i)}]^2 \right\}$$

Experimental squeezing criterion

$$(\delta p_{12})_{sq}^2 < \frac{1}{4} \eta \langle \hat{N} \rangle = \eta (\delta p_1)_{proj}^2$$



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Near spin-squeezing

Experimental parameters

Projection noise measurement

$$(\delta p_1)_{proj}^2 = \frac{1}{k-1} \sum_{i=1}^K [p_{p,1}^{(i)}]^2 \propto \frac{1}{4} \langle \hat{N} \rangle$$

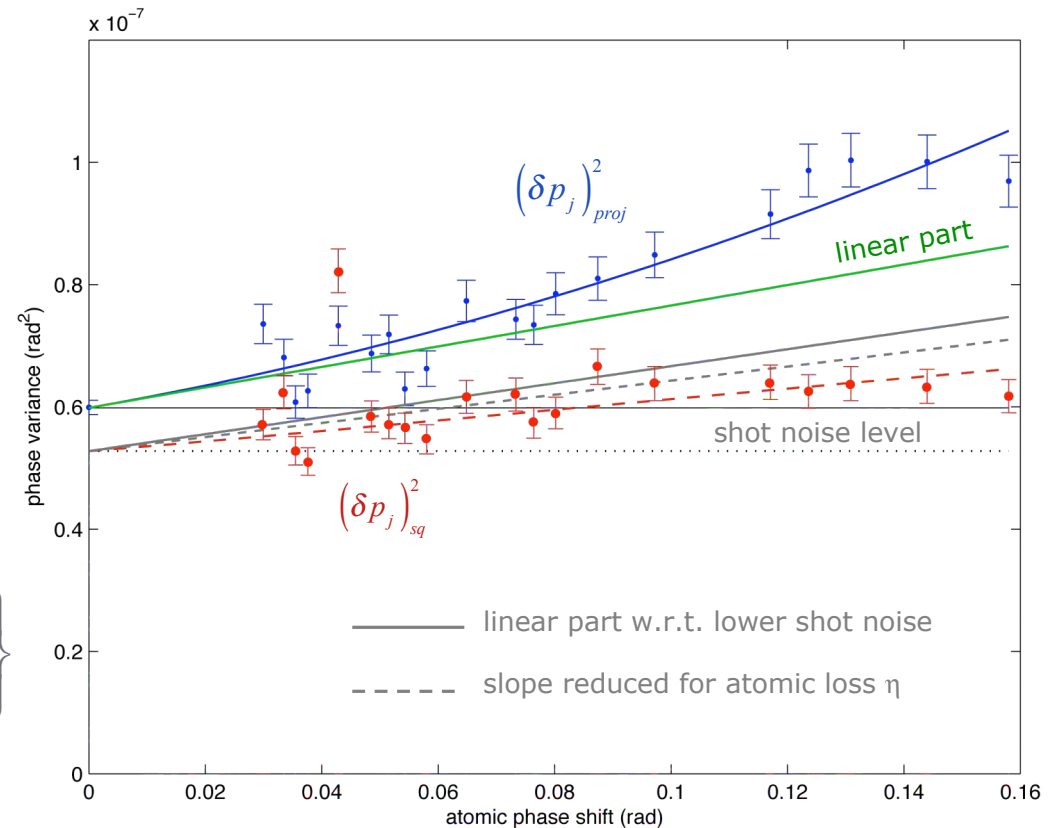
Squeezing measurement - in principle

$$(\delta p_j)_{sq}^2 = \frac{1}{2k-1} \sum_{i=1}^K [p_{p,j}^{(i)} - \chi_{sq} p_{p,j+1}^{(i)}]^2$$

$$\chi_{sq} = \min_{\chi} \left\{ \frac{1}{2k-1} \sum_{i=1}^K [p_{p,j}^{(i)} - \chi p_{p,j+1}^{(i)}]^2 \right\}$$

Experimental squeezing criterion

$$(\delta p_{12})_{sq}^2 < \frac{1}{4} \eta \langle \hat{N} \rangle = \eta (\delta p_1)_{proj}^2$$



Pulse durations: 4 μ s

Pulse separation: 16 μ s

Probe power: 10 μ W \sim 1.6 10^8 photons

Reference power: 50 μ W \sim 8 10^8 photons



Near spin-squeezing

Experimental parameters

Projection noise measurement

$$\left(\delta p_1\right)_{proj}^2 = \frac{1}{k-1} \sum_{i=1}^K \left[p_{p,1}^{(i)} \right]^2 \propto \frac{1}{4} \langle \hat{N} \rangle$$

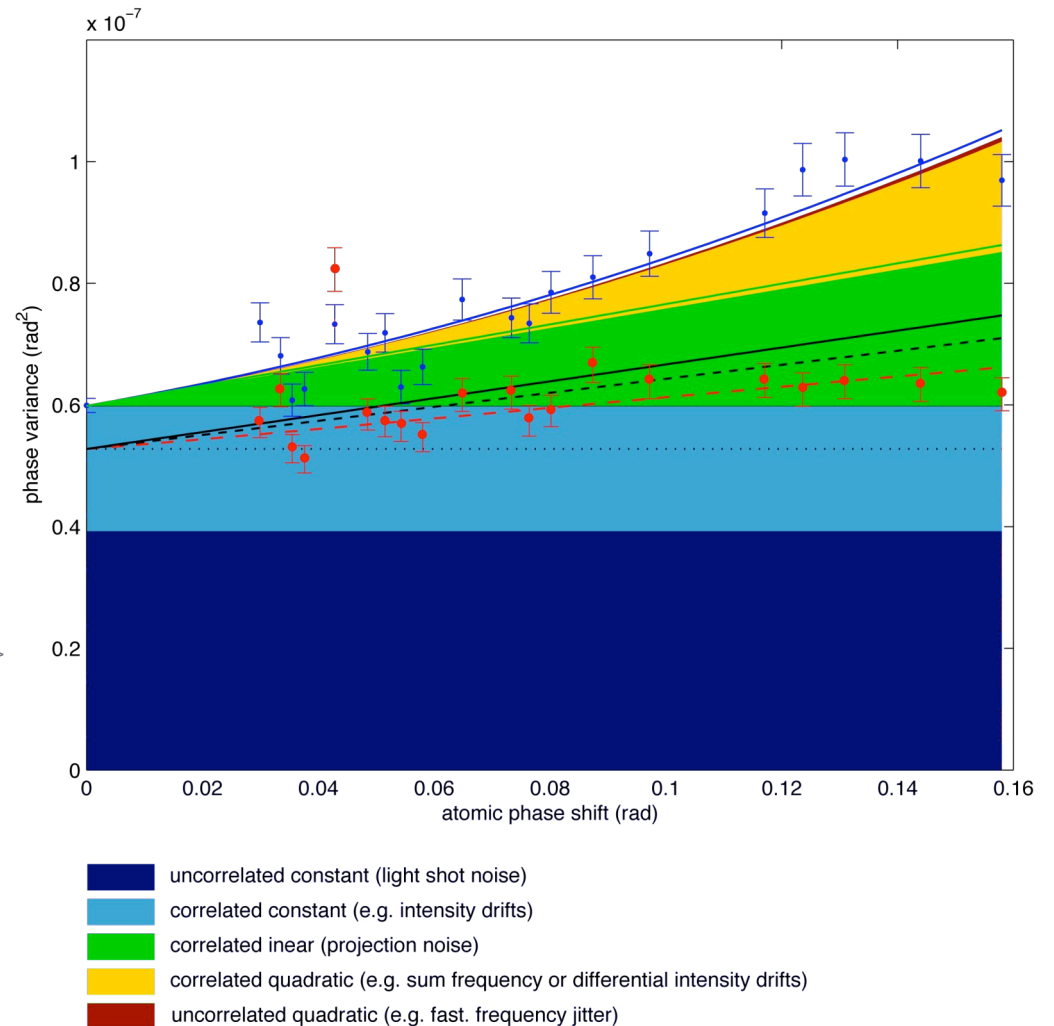
Squeezing measurement - in principle

$$\left(\delta p_j\right)_{sq}^2 = \frac{1}{2k-1} \sum_{i=1}^K \left[p_{p,j}^{(i)} - \chi_{sq} p_{p,j+1}^{(i)} \right]^2$$

$$\chi_{sq} = \min_{\chi} \left\{ \frac{1}{2k-1} \sum_{i=1}^K \left[p_{p,j}^{(i)} - \chi p_{p,j+1}^{(i)} \right]^2 \right\}$$

Experimental squeezing criterion

$$\left(\delta p_{12}\right)_{sq}^2 < \frac{1}{4} \eta \langle \hat{N} \rangle = \eta \left(\delta p_1\right)_{proj}^2$$





Conclusion

Shot noise limited phase-sensitive measurement

Non destructive measurements of quantum state dynamics

Light shift

Observation of the reduction of quantum projection noise

Limited by classical background noise

QND measurement



Outlook

Two input port interferometer setup¹

Sensitivity to difference frequency of 2 colours -> narrower linewidth

Reduced sensitivity to acoustic noise -> reference pulses superfluous

No more - that should work



Acknowledgements



The Danish National Research Foundation
QUANTOP

European Union
QAUAC
QOVAQIAL
QUICOV
QWAP





And thanks for your attention

And sorry for speaking too long...

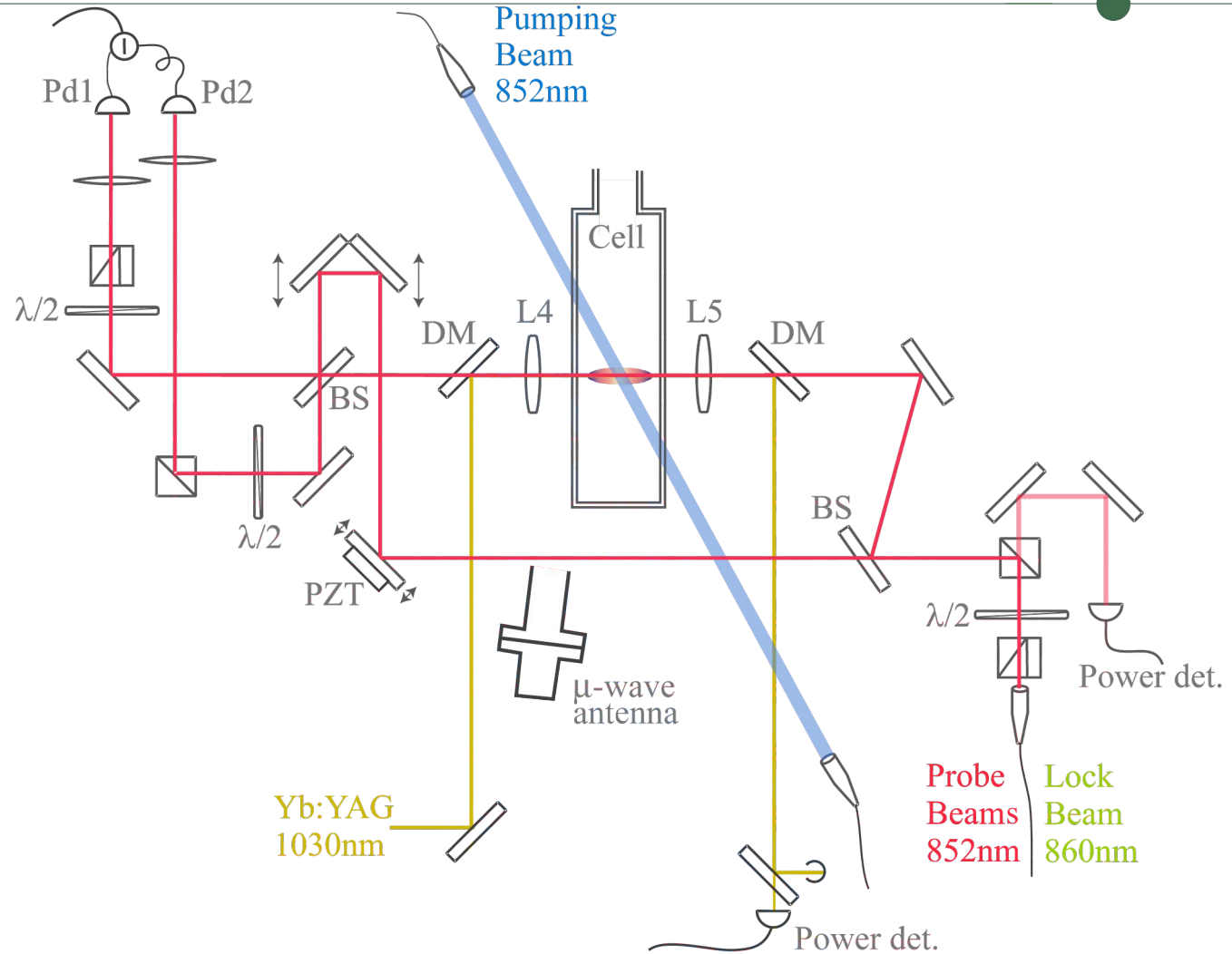


Any questions?



Two input interferometer

Benefits:



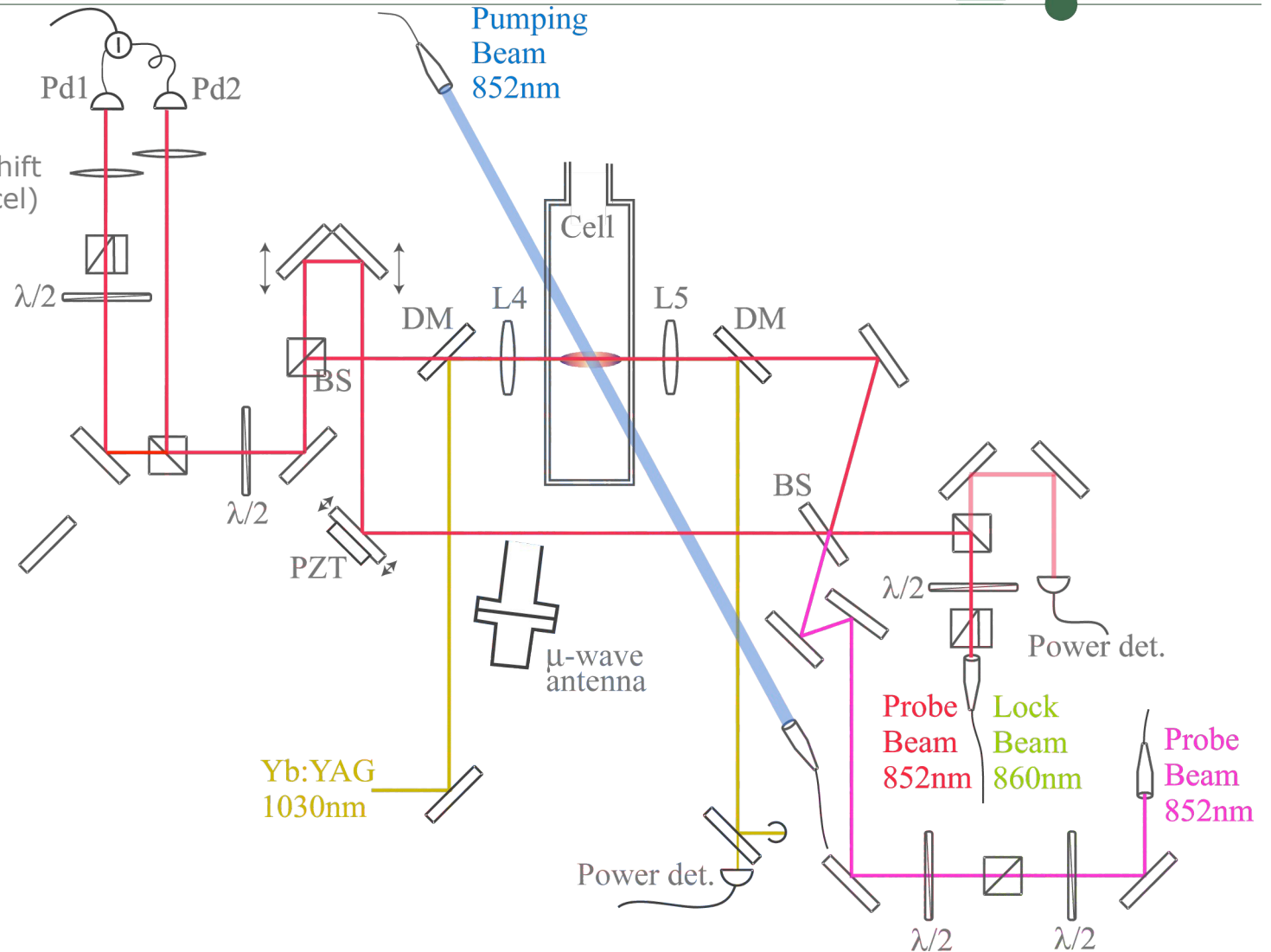


Two input interferometer

Benefits:

Eliminates light shift
(two colours cancel)

Insensitive to
mechanical and
acoustic noise





Noise sources in atomic measurements

Quantum and classical noise sources

Noise in output signal

$$\hat{n}_-^{out} = \hat{n}_-^{in} + 4\tilde{\kappa}\tau \langle \hat{n} \rangle \hat{J}_z$$

$$(\delta \hat{n}_-^{out})^2 = (\delta \hat{n}_-^{in})^2 + 4\tilde{\kappa}\tau \langle \hat{n} \rangle (\delta \hat{J}_z)^2$$

$$= \langle \hat{n} \rangle + \tilde{\kappa}\tau \langle \hat{n} \rangle \langle \hat{N} \rangle + K \langle \hat{N} \rangle^2 + E$$



light
shot



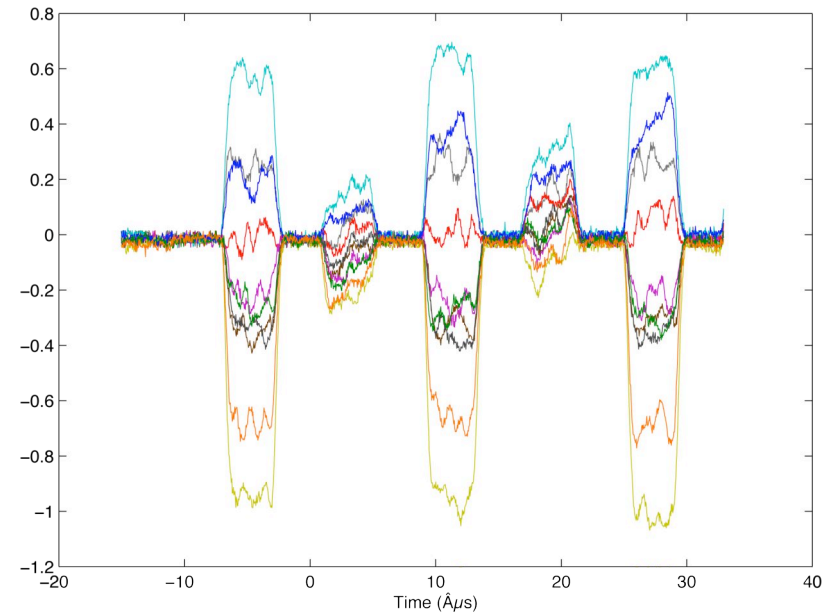
atomic
projection



classical
garbage



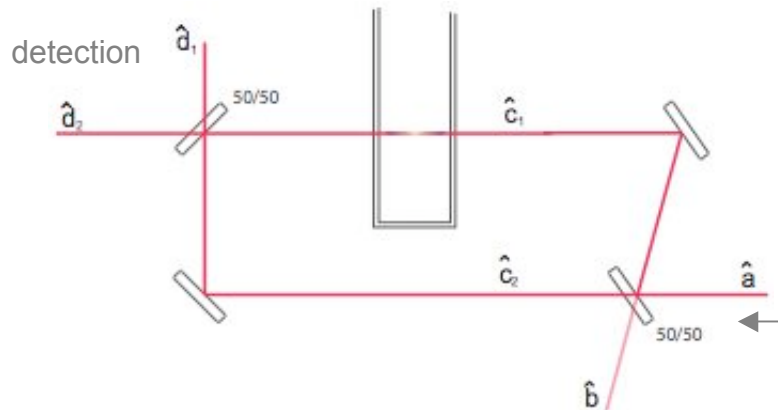
electr.





X Interferometric measurement

Interaction cast in continuous light variables



Schwinger representation

$$\hat{S}_{kx} = \frac{1}{2}(\hat{a}_k^\dagger \hat{b}_k + \hat{b}_k^\dagger \hat{a}_k)$$

$$\hat{S}_{ky} = \frac{-i}{2}(\hat{a}_k^\dagger \hat{b}_k - \hat{b}_k^\dagger \hat{a}_k)$$

$$\hat{S}_{kz} = \frac{1}{2}(\hat{a}_k^\dagger \hat{a}_k - \hat{b}_k^\dagger \hat{b}_k)$$

$$\hat{n}_k = \hat{a}_k^\dagger \hat{a}_k + \hat{b}_k^\dagger \hat{b}_k$$

$$\langle \hat{a}_k^\dagger \hat{a}_k \rangle = \langle \hat{n}_k \rangle$$

$$\langle \hat{b}_k^\dagger \hat{b}_k \rangle = 0$$

$$\langle \hat{S}_{kx}^c \rangle = \langle \hat{S}_{kz}^c \rangle = 0$$

$$\langle \hat{S}_{ky}^c \rangle = \frac{1}{2} \langle \hat{n}_k \rangle$$

Interaction Hamiltonian

$$\hat{H}_{k \text{ int}} = \hbar 4 \tilde{\kappa}_k \left(\frac{1}{2} \hat{n}_k + \hat{S}_{kz}^c \right) \hat{N}_k \Rightarrow$$

$$\hat{H}_{\text{int}} = \hbar 4 \left[\tilde{\kappa}_3 \left(\frac{1}{2} \hat{n}_3 + \hat{S}_{3z}^c \right) \hat{N}_3 + \tilde{\kappa}_4 \left(\frac{1}{2} \hat{n}_4 + \hat{S}_{4z}^c \right) \hat{N}_4 \right]$$

Operator input-output relations

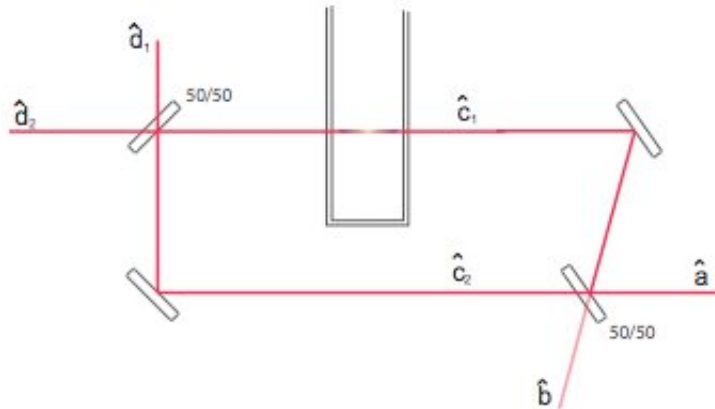
$$\hat{S}_{kz}^{c, \text{out}} = \hat{S}_{kz}^{c, \text{in}} \quad , \quad \hat{S}_{kx}^{c, \text{out}} = \hat{S}_{kx}^{c, \text{in}} + 2 \tilde{\kappa}_k \tau \langle \hat{n}_k \rangle \hat{N}_k$$

$$\hat{J}_z^{c, \text{out}} = \hat{J}_z^{c, \text{in}} \quad , \quad \hat{J}_y^{\text{out}} = \hat{J}_y^{\text{in}} + 2 \tilde{\kappa}_k \tau \hat{N}_k \left(\frac{1}{2} \langle \hat{n}_k \rangle + \hat{S}_{kz}^c \right)$$



X Interferometric measurement

Quantum picture



$$i_- \propto \langle \hat{S}_{3z}^d + \hat{S}_{4z}^d \rangle = \hat{S}_{3x}^{c,out} + \hat{S}_{4x}^{c,out} = \hat{S}_{kx}^{c,in} + 2\tilde{\kappa}_k \tau \langle \hat{n}_k \rangle \hat{J}_z$$

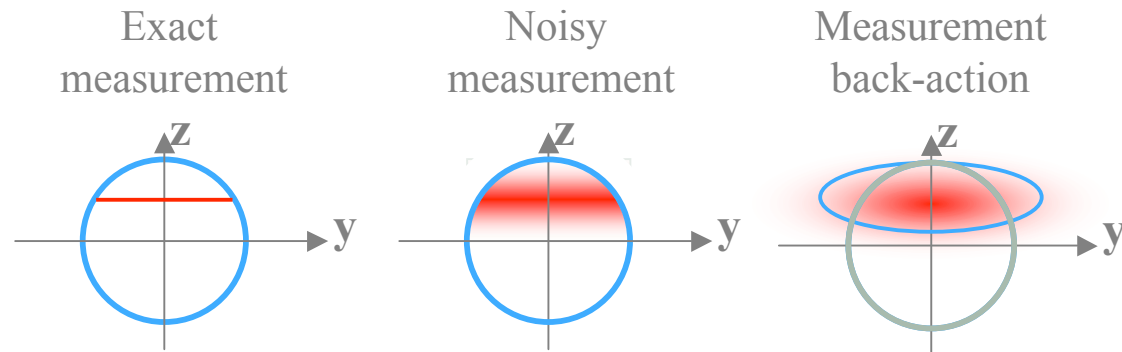
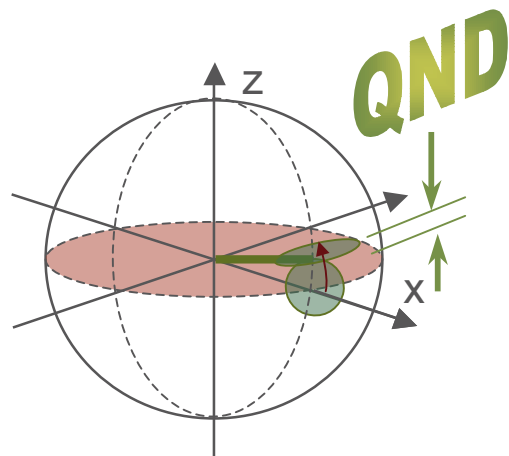
$$\hat{S}_{kz}^{c,out} = \hat{S}_{kz}^{c,in} \quad , \quad \hat{S}_{kx}^{c,out} = \hat{S}_{kx}^{c,in} + 2\tilde{\kappa}_k \tau \langle \hat{n}_k \rangle \hat{J}_z$$

$$\hat{J}_z^{c,out} = \hat{J}_z^{c,in} \quad , \quad \hat{J}_y^{out} = \hat{J}_y^{in} + 2\tilde{\kappa}_k \tau \hat{N} \left(\frac{1}{2} \langle \hat{n}_k \rangle + \hat{S}_{kz}^c \right)$$



X Quantum non-demolition measurement

Squeezing by interaction and subsequent measurement of light pulse



QND interaction Hamiltonian

$$\hat{H}_{\text{int}} = \hbar 4\tilde{\kappa} \left(\frac{\hat{n}}{2} + \hat{S}_z^c \right) \hat{J}_z$$

\Rightarrow

Operator input-output relations

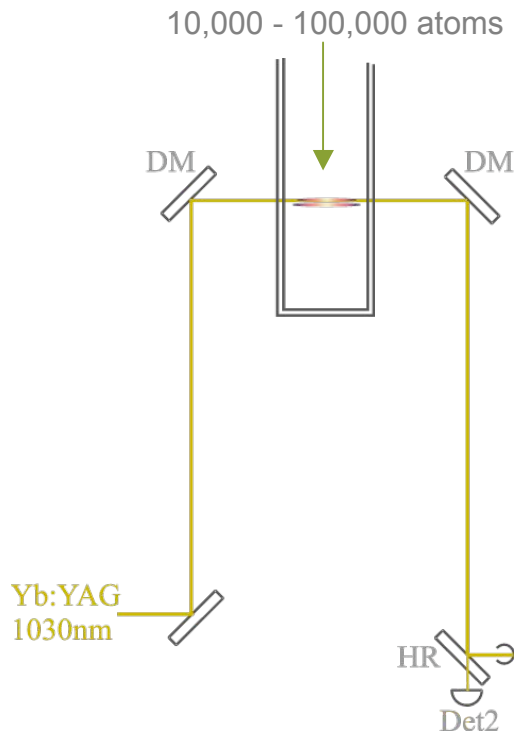
$$\begin{aligned} \hat{S}_z^{c,out} &= \hat{S}_z^{c,in} , & \hat{S}_x^{c,out} &= \hat{S}_x^{c,in} + 2\tilde{\kappa}\tau \langle \hat{n} \rangle \hat{J}_z \\ \hat{J}_z^{c,out} &= \hat{J}_z^{c,in} , & \hat{J}_y^{out} &= \hat{J}_y^{in} + 2\tilde{\kappa}\tau \hat{N} \left(\frac{\langle \hat{n} \rangle}{2} + \hat{S}_z^c \right) \end{aligned}$$



Far off-resonant trap (FORT)

Setup →

Atoms trapped in MOT
→ sub-doppler cooled
→ loaded into FORT



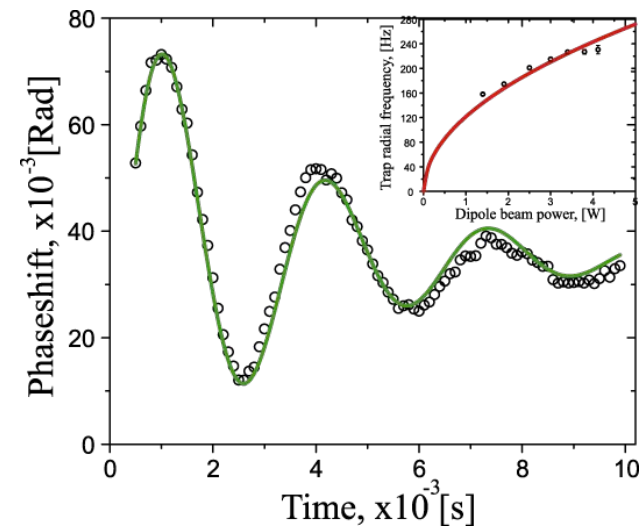
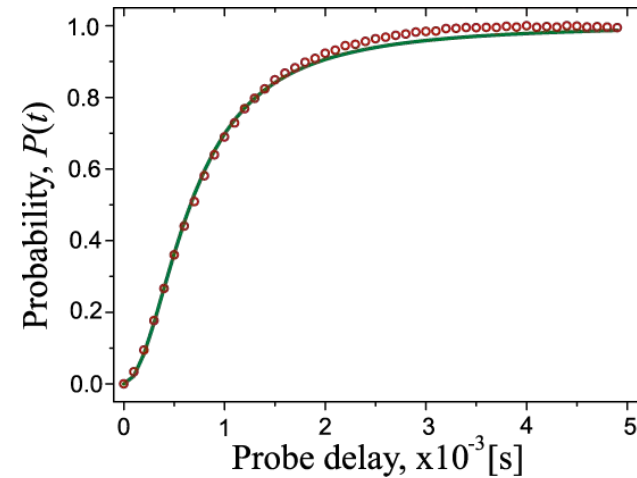
Expansion rate
→ temperature

$T \approx 15\mu\text{K}$

Only 1 probe colour!

Atom breathing
→ trap radial
oscillation frequency

$f = 180 - 220\text{Hz}$

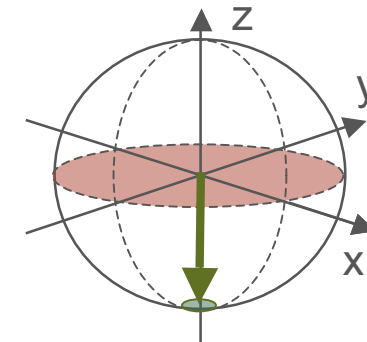
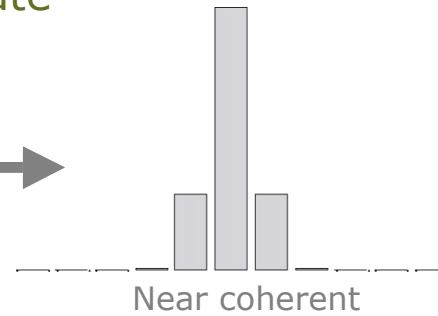




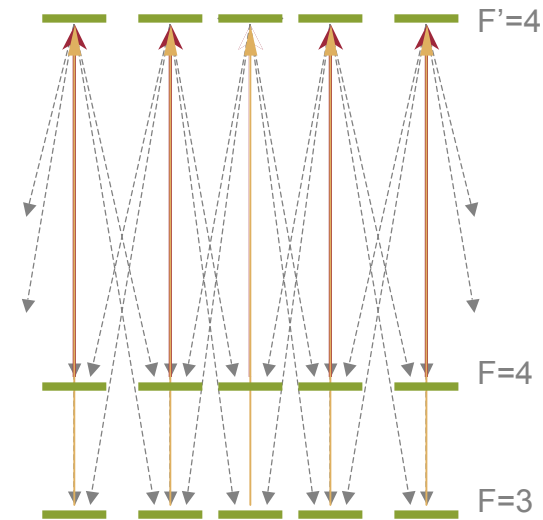
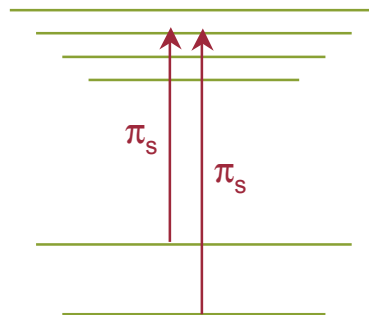
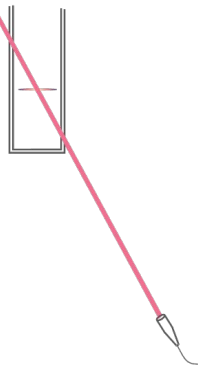
Optical pumping to the $m_f=0$ sublevel

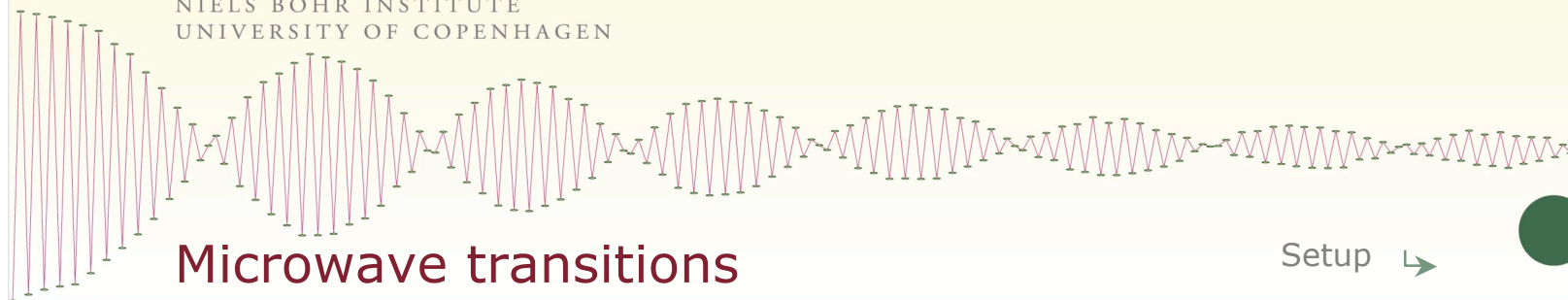
Setup \rightarrow

Generating a coherent state of a 2-level system



Pumping Beam
852nm





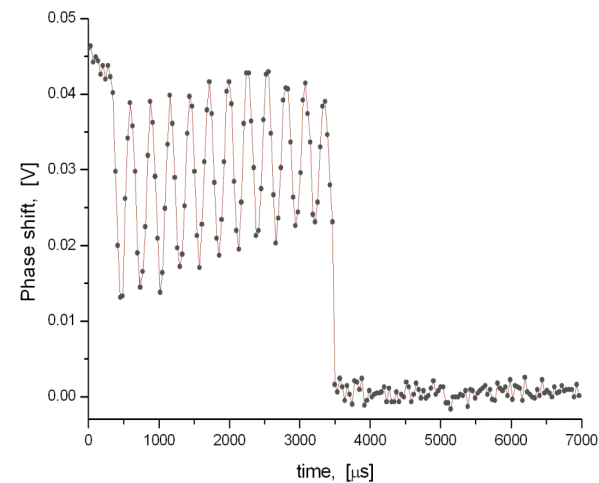
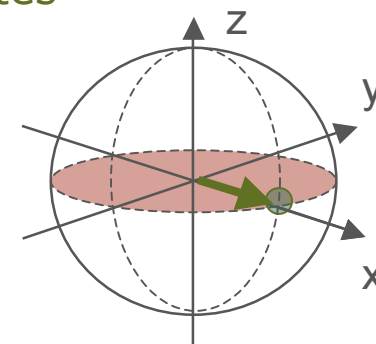
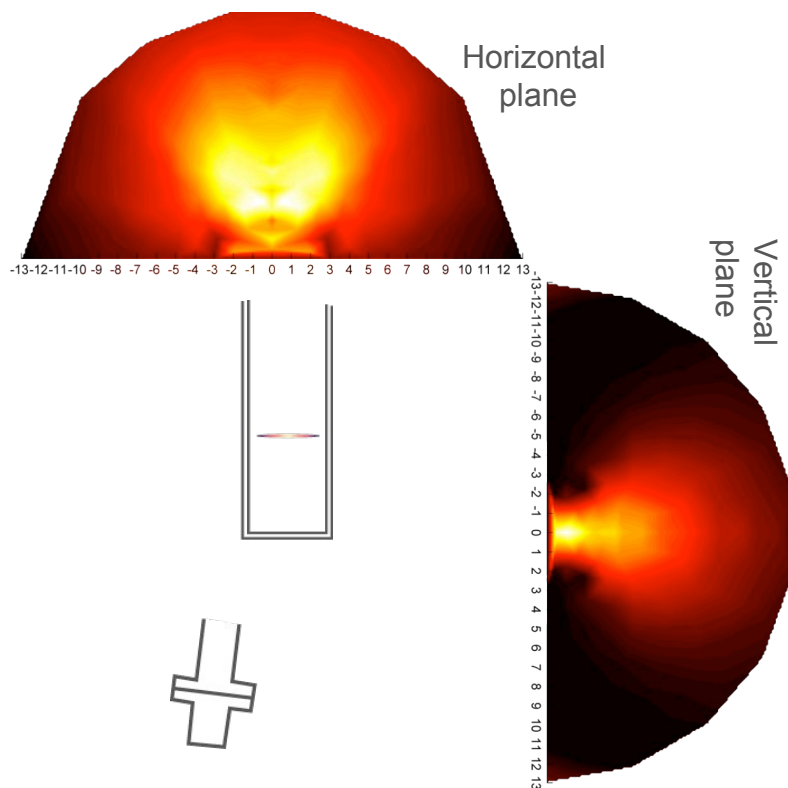
Microwave transitions

Setup →



Magnetic dipole transitions btw hyperfine ground states

Sawed-off waveguide



Optical pumping efficiency:
fraction of atoms in Rabi oscillations,